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LETTER TO THE EDITOR

A CrO₂-based magnetic tunnel junctionA Barry[†], J M D Coey[†] and M Viret[‡][†] Department of Physics, Trinity College, Dublin 2, Ireland[‡] Service de Physique de l'Etat Condensée, Orme des Merisiers, CEA-Saclay, 91191 Gif-sur-Yvette, France

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Abstract. A tunnel junction based on the half-metallic oxide CrO₂ uses a native oxide barrier layer and a cobalt top electrode. The $I:V$ characteristic is fitted to the Simmons model with $\phi = 0.76$ eV and $t = 2.0$ nm. The magnetoresistance is positive with $\Delta R/R = 1.0\%$ at 77 K.

Half-metallic oxides such as CrO₂ are considered ideal materials for spin electronic devices on account of their spin-split d band and the absence of sp conduction electrons, giving a single sign of spin polarization at the Fermi level [1]. In a recent study of the spin polarization using Andreev reflection in the liquid helium range, CrO₂ was found to exhibit the greatest value of spin polarization (90%) of any of the seven 3d metals examined [2]. The oxide has the rutile structure with a 3d² ($t_{2g}^2 \uparrow$) configuration and a magnetic moment of 2.0 μ_B /formula unit at 4.2 K. Here we present first results on a spin-polarized junction based on the same CrO₂ films as were used for the Andreev reflection measurements.

The films were deposited on a (110)-oriented TiO₂ substrate by the thermal decomposition of CrO₃, as described previously [3]. The film thickness was ~ 800 nm. Full details of their characterization and properties are found in a recent thesis [4].

For the present study, UV lithography was used to pattern the CrO₂ films. The films were wet etched to produce a bar of width 5 μm using an industrial chrome etch composed of ammonium ceric nitrate and nitric acid. Etch rates were in the range 0.008–0.011 $\mu\text{m s}^{-1}$, depending on film thickness and quality. The tunnel barrier was a native oxide layer and was grown by exposing the film to air for several weeks. As no attempt was made to reduce the CrO₂ films after deposition, the barrier is probably composed of antiferromagnetic Cr₂O₃, the equilibrium chromium–oxygen phase in ambient conditions. Our experience with the reduction of CrO₂ at elevated temperatures has shown that after the initial reduction, the thickness of the antiferromagnetic layer does not depend on time; it acts as a passivation layer. A Co top electrode, 150 nm in thickness, was deposited by RF sputtering, patterned using UV lithography and etched by ion beam milling. A schematic drawing of the device with dimensions $5 \times 5 \mu\text{m}^2$ is shown in figure 1. Contacts were made to copper pads by wire bonding using Al wire 50 μm thick.

The $I:V$ curve is illustrated in figure 2. It shows a non-linear characteristic consistent with tunnelling across a barrier with a parabolic variation of the conductance (dI/dV) versus voltage (V). The junction resistance was ~ 60 k Ω . A fit to the Simmons model [5] yields barrier thicknesses and barrier heights of 2.0 nm and 0.76 eV at 13 K and 0.79 eV and 2.5 nm at 77 K.

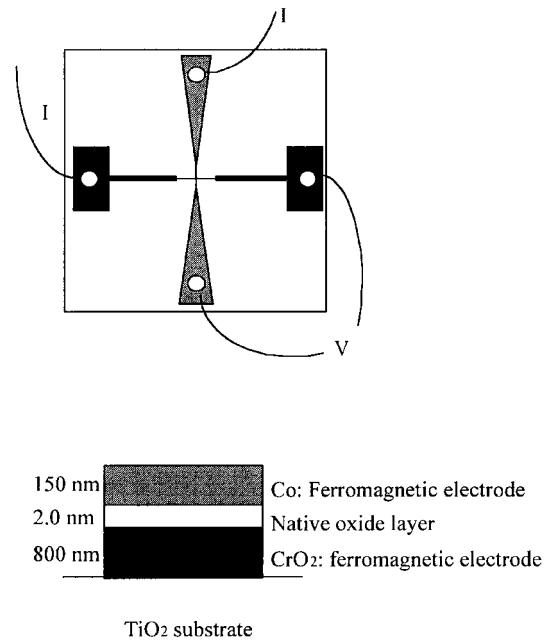


Figure 1. The ferromagnetic tunnel junction: CrO₂/I/Co in the cross-stripe geometry with the current perpendicular to the plane (CPP). Dimensions are in nm.

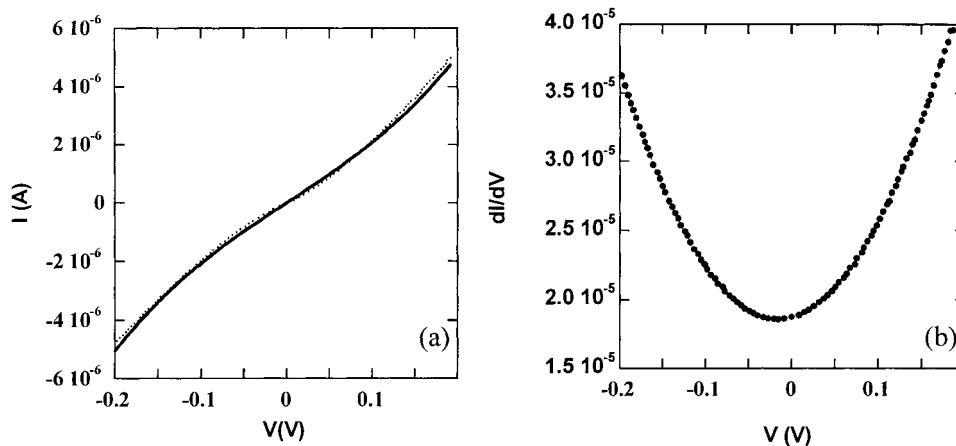


Figure 2. (a) I : V characteristics at $T = 13$ K for a tunnel junction of the type CrO₂/barrier layer/Co fitted to the Simmons model (solid black line) and (b) dI/dV versus V for the same junction.

The R : H curve measured with the field parallel to the $5 \mu\text{m}$ CrO₂ line is shown in figure 3. The switching field of the CrO₂ electrode is 20 mT and that of the cobalt electrode is 50 mT. For intermediate fields they adopt an antiparallel configuration with a resistance 1% higher than that for the parallel configuration.

The magnitude of the spin-polarized tunnel magnetoresistance observed at 77 K is quite

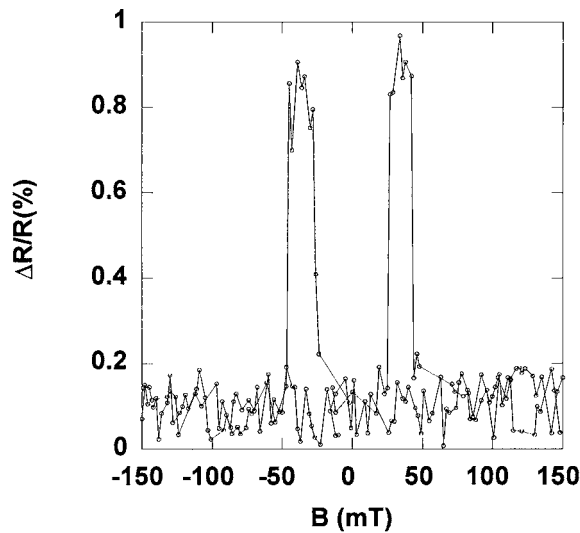


Figure 3. The magnetoresistance curve at 77 K of the $25 \mu\text{m}^2$ $\text{CrO}_2/\text{I}/\text{Co}$ tunnel junction.

small. Higher values up to 7% at 77 K and 52% at 4.2 K were found for pressed powder compacts [6, 7] and partly oxidized films [8]. The MR effect is considerably smaller than that expected for complete spin polarization and it is strongly temperature dependent. Part of the problem may be that the barrier oxide Cr_2O_3 is antiferromagnetic, with an antiferromagnetic axis that will bear no relation to the ferromagnetic directions in the adjacent Co or CrO_2 layers. Tunnelling electrons are then subject to a fluctuating exchange field that could induce spin depolarization [9]. Antiferromagnetic magnons can be excited in the barrier. Future work on spin-dependent tunnelling in CrO_2 -based structures should investigate non-magnetic barriers such as TiO_2 and ZrO_2 .

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